

# **Complete Solutions Manual to Accompany**

---

## **A First Course in Differential Equations with Modeling Applications**

**ELEVENTH EDITION, METRIC VERSION**

**And**

## **Differential Equations with Boundary-Value Problems**

**NINTH EDITION, METRIC VERSION**

**Dennis G. Zill**

Loyola Marymount University,  
Los Angeles, CA

Prepared by

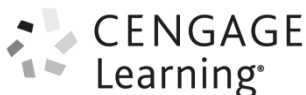
**Roberto Martinez**

Loyola Marymount University, Los Angeles, CA

Metric Version Prepared by

**Aly El-Iraki**

Professor Emeritus, Alexandria University, Egypt



---

Australia • Brazil • Mexico • Singapore • United Kingdom • United States



© 2018 Cengage Learning

ALL RIGHTS RESERVED. No part of this work covered by the copyright herein may be reproduced, transmitted, stored, or used in any form or by any means graphic, electronic, or mechanical, including but not limited to photocopying, recording, scanning, digitizing, taping, Web distribution, information networks, or information storage and retrieval systems, except as permitted under Section 107 or 108 of the 1976 United States Copyright Act, without the prior written permission of the publisher except as may be permitted by the license terms below.

For product information and technology assistance, contact us at **Cengage Learning Customer & Sales Support, 1-800-354-9706**.

For permission to use material from this text or product, submit all requests online at **[www.cengage.com/permissions](http://www.cengage.com/permissions)** Further permissions questions can be emailed to **[permissionrequest@cengage.com](mailto:permissionrequest@cengage.com)**.

Ninth Edition, Metric Version  
ISBN-13: 978-1-337-55658-3  
ISBN-10: 1-337-55658-0

Eleventh Edition, Metric Version  
ISBN-13: 978-1-337-55666-8  
ISBN-10: 1-337-55666-1

**Cengage Learning**  
200 First Stamford Place, 4th Floor  
Stamford, CT 06902  
USA

Cengage Learning is a leading provider of customized learning solutions with office locations around the globe, including Singapore, the United Kingdom, Australia, Mexico, Brazil, and Japan. Locate your local office at: **[www.cengage.com/global](http://www.cengage.com/global)**.

Cengage Learning products are represented in Canada by Nelson Education, Ltd.

To learn more about Cengage Learning Solutions, visit **[www.cengage.com](http://www.cengage.com)**.

Purchase any of our products at your local college store or at our preferred online store **[www.cengagebrain.com](http://www.cengagebrain.com)**.

**NOTE: UNDER NO CIRCUMSTANCES MAY THIS MATERIAL OR ANY PORTION THEREOF BE SOLD, LICENSED, AUCTIONED, OR OTHERWISE REDISTRIBUTED EXCEPT AS MAY BE PERMITTED BY THE LICENSE TERMS HEREIN.**

**READ IMPORTANT LICENSE INFORMATION**

Dear Professor or Other Supplement Recipient:

Cengage Learning has provided you with this product (the “Supplement”) for your review and, to the extent that you adopt the associated textbook for use in connection with your course (the “Course”), you and your students who purchase the textbook may use the Supplement as described below. Cengage Learning has established these use limitations in response to concerns raised by authors, professors, and other users regarding the pedagogical problems stemming from unlimited distribution of Supplements.

Cengage Learning hereby grants you a nontransferable license to use the Supplement in connection with the Course, subject to the following conditions. The Supplement is for your personal, noncommercial use only and may not be reproduced, posted electronically or distributed, except that portions of the Supplement may be provided to your students **IN PRINT FORM ONLY** in connection with your instruction of the Course, so long as such students are advised that they may not copy or distribute any portion of the Supplement to any

third party. You may not sell, license, auction, or otherwise re-distribute the Supplement in any form. We ask that you take reasonable steps to protect the Supplement from unauthorized use, reproduction, or distribution. Your use of the Supplement indicates your acceptance of the conditions set forth in this Agreement. If you do not accept these conditions, you must return the Supplement unused within 30 days of receipt.

All rights (including without limitation, copyrights, patents, and trade secrets) in the Supplement are and will remain the sole and exclusive property of Cengage Learning and/or its licensors. The Supplement is furnished by Cengage Learning on an “as is” basis without any warranties, express or implied. This Agreement will be governed by and construed pursuant to the laws of the State of New York, without regard to such State’s conflict of law rules.

Thank you for your assistance in helping to safeguard the integrity of the content contained in this Supplement. We trust you find the Supplement a useful teaching tool.

# Chapter 1

## Introduction to Differential Equations

1.1

Definitions and Terminology

1. Second order; linear
2. Third order; nonlinear because of  $(dy/dx)^4$
3. Fourth order; linear
4. Second order; nonlinear because of  $\cos(r + u)$
5. Second order; nonlinear because of  $(dy/dx)^2$  or  $\sqrt{1 + (dy/dx)^2}$
6. Second order; nonlinear because of  $R^2$
7. Third order; linear
8. Second order; nonlinear because of  $\dot{x}^2$
9. Writing the differential equation in the form  $x(dy/dx) + y^2 = 1$ , we see that it is nonlinear in  $y$  because of  $y^2$ . However, writing it in the form  $(y^2 - 1)(dx/dy) + x = 0$ , we see that it is linear in  $x$ .
10. Writing the differential equation in the form  $u(dv/du) + (1 + u)v = ue^u$  we see that it is linear in  $v$ . However, writing it in the form  $(v + uv - ue^u)(du/dv) + u = 0$ , we see that it is nonlinear in  $u$ .
11. From  $y = e^{-x/2}$  we obtain  $y' = -\frac{1}{2}e^{-x/2}$ . Then  $2y' + y = -e^{-x/2} + e^{-x/2} = 0$ .
12. From  $y = \frac{6}{5} - \frac{6}{5}e^{-20t}$  we obtain  $dy/dt = 24e^{-20t}$ , so that
$$\frac{dy}{dt} + 20y = 24e^{-20t} + 20\left(\frac{6}{5} - \frac{6}{5}e^{-20t}\right) = 24.$$
13. From  $y = e^{3x} \cos 2x$  we obtain  $y' = 3e^{3x} \cos 2x - 2e^{3x} \sin 2x$  and  $y'' = 5e^{3x} \cos 2x - 12e^{3x} \sin 2x$ , so that  $y'' - 6y' + 13y = 0$ .

58. Writing  $\cos^2 x = \frac{1}{2}(1 + \cos 2x)$  and applying  $D(D^2 + 4)$  to the differential equation we obtain

$$D(D^2 + 4)(D^2 + 4) = D(D^2 + 4)^2 = 0.$$

Then

$$y = \underbrace{c_1 \cos 2x + c_2 \sin 2x}_{y_c} + c_3 x \cos 2x + c_4 x \sin 2x + c_5$$

and  $y_p = Ax \cos 2x + Bx \sin 2x + C$ . Substituting  $y_p$  into the differential equation yields

$$-4A \sin 2x + 4B \cos 2x + 4C = \frac{1}{2} + \frac{1}{2} \cos 2x.$$

Equating coefficients gives  $A = 0$ ,  $B = 1/8$ , and  $C = 1/8$ . The general solution is

$$y = c_1 \cos 2x + c_2 \sin 2x + \frac{1}{8}x \sin 2x + \frac{1}{8}.$$

59. Applying  $D^3$  to the differential equation we obtain

$$D^3(D^3 + 8D^2) = D^5(D + 8) = 0.$$

Then

$$y = \underbrace{c_1 + c_2 x + c_3 e^{-8x}}_{y_c} + c_4 x^2 + c_5 x^3 + c_6 x^4$$

and  $y_p = Ax^2 + Bx^3 + Cx^4$ . Substituting  $y_p$  into the differential equation yields

$$16A + 6B + (48B + 24C)x + 96Cx^2 = 2 + 9x - 6x^2.$$

Equating coefficients gives

$$16A + 6B = 2$$

$$48B + 24C = 9$$

$$96C = -6.$$

Then  $A = 11/256$ ,  $B = 7/32$ , and  $C = -1/16$ , and the general solution is

$$y = c_1 + c_2 x + c_3 e^{-8x} + \frac{11}{256}x^2 + \frac{7}{32}x^3 - \frac{1}{16}x^4.$$

60. Applying  $D(D - 1)^2(D + 1)$  to the differential equation we obtain

$$D(D - 1)^2(D + 1)(D^3 - D^2 + D - 1) = D(D - 1)^3(D + 1)(D^2 + 1) = 0.$$

Then

$$y = \underbrace{c_1 e^x + c_2 \cos x + c_3 \sin x}_{y_c} + c_4 + c_5 e^{-x} + c_6 x e^x + c_7 x^2 e^x$$

57. Solving  $\frac{1}{2}q'' + 10q' + 100q = 150$  we obtain  $q(t) = e^{-10t}(c_1 \cos 10t + c_2 \sin 10t) + 3/2$ . The initial conditions  $q(0) = 1$  and  $q'(0) = 0$  imply  $c_1 = c_2 = -1/2$ . Thus

$$q(t) = -\frac{1}{2}e^{-10t}(\cos 10t + \sin 10t) + \frac{3}{2}.$$

As  $t \rightarrow \infty$ ,  $q(t) \rightarrow 3/2$ .

58. In Problem 54 it is shown that the amplitude of the steady-state current is  $E_0/Z$ , where  $Z = \sqrt{X^2 + R^2}$  and  $X = L\gamma - 1/C\gamma$ . Since  $E_0$  is constant the amplitude will be a maximum when  $Z$  is a minimum. Since  $R$  is constant,  $Z$  will be a minimum when  $X = 0$ . Solving  $L\gamma - 1/C\gamma = 0$  for  $\gamma$  we obtain  $\gamma = 1/\sqrt{LC}$ . The maximum amplitude will be  $E_0/R$ .

59. By Problem 54 the amplitude of the steady-state current is  $E_0/Z$ , where  $Z = \sqrt{X^2 + R^2}$  and  $X = L\gamma - 1/C\gamma$ . Since  $E_0$  is constant the amplitude will be a maximum when  $Z$  is a minimum. Since  $R$  is constant,  $Z$  will be a minimum when  $X = 0$ . Solving  $L\gamma - 1/C\gamma = 0$  for  $C$  we obtain  $C = 1/L\gamma^2$ .

60. Solving  $0.1q'' + 10q = 100 \sin \gamma t$  we obtain

$$q(t) = c_1 \cos 10t + c_2 \sin 10t + q_p(t)$$

where  $q_p(t) = A \sin \gamma t + B \cos \gamma t$ . Substituting  $q_p(t)$  into the differential equation we find

$$(100 - \gamma^2)A \sin \gamma t + (100 - \gamma^2)B \cos \gamma t = 100 \sin \gamma t.$$

Equating coefficients we obtain  $A = 100/(100 - \gamma^2)$  and  $B = 0$ . Thus,  $q_p(t) = \frac{100}{100 - \gamma^2} \sin \gamma t$ . The initial conditions  $q(0) = q'(0) = 0$  imply  $c_1 = 0$  and  $c_2 = -10\gamma/(100 - \gamma^2)$ . The charge is

$$q(t) = \frac{10}{100 - \gamma^2}(10 \sin \gamma t - \gamma \sin 10t)$$

and the current is

$$i(t) = \frac{100\gamma}{100 - \gamma^2}(\cos \gamma t - \cos 10t).$$

61. In an  $LC$ -series circuit there is no resistor, so the differential equation is

$$L \frac{d^2 q}{dt^2} + \frac{1}{C} q = E(t).$$

Then  $q(t) = c_1 \cos(t/\sqrt{LC}) + c_2 \sin(t/\sqrt{LC}) + q_p(t)$  where  $q_p(t) = A \sin \gamma t + B \cos \gamma t$ .

Substituting  $q_p(t)$  into the differential equation we find

$$\left(\frac{1}{C} - L\gamma^2\right) A \sin \gamma t + \left(\frac{1}{C} - L\gamma^2\right) B \cos \gamma t = E_0 \cos \gamma t.$$

Equating coefficients we obtain  $A = 0$  and  $B = E_0 C/(1 - LC\gamma^2)$ . Thus, the charge is

$$q(t) = c_1 \cos \frac{1}{\sqrt{LC}} t + c_2 \sin \frac{1}{\sqrt{LC}} t + \frac{E_0 C}{1 - LC\gamma^2} \cos \gamma t.$$

From the last result and using  $\nu = 3/2$  we obtain

$$\begin{aligned} 3J_{3/2}(x) &= xJ_{5/2}(x) + xJ_{1/2}(x) \\ J_{5/2}(x) &= \frac{3}{x}J_{3/2}(x) - J_{1/2}(x) \\ &= \frac{3}{x}\sqrt{\frac{2}{\pi x}}\left(\frac{\sin x}{x} - \cos x\right) - \sqrt{\frac{2}{\pi x}}\sin x \\ &= \sqrt{\frac{2}{\pi x}}\left[\left(\frac{3}{x^2} - 1\right)\sin x - \frac{3\cos x}{x}\right] \end{aligned}$$

From the last result and using  $\nu = 5/2$  we obtain

$$\begin{aligned} 5J_{5/2}(x) &= xJ_{7/2}(x) + xJ_{3/2}(x) \\ J_{7/2}(x) &= \frac{5}{x}J_{5/2}(x) - J_{3/2}(x) \\ &= \frac{5}{x}\sqrt{\frac{2}{\pi x}}\left(\frac{3\sin x}{x^2} - \frac{3\cos x}{x} - \sin x\right) - \sqrt{\frac{2}{\pi x}}\left(\frac{\sin x}{x} - \cos x\right) \\ &= \sqrt{\frac{2}{\pi x}}\left[\left(\frac{15}{x^3} - \frac{6}{x}\right)\sin x - \left(\frac{15}{x^2} - 1\right)\cos x\right] \end{aligned}$$

**33. (a)** To find the spherical Bessel functions  $j_1(x)$  and  $j_2(x)$  we use the first formula in (30),

$$j_n(x) = \sqrt{\frac{\pi}{2x}} J_{n+1/2}(x)$$

with  $n = 1$  and  $n = 2$ ,

$$j_1(x) = \sqrt{\frac{\pi}{2x}} J_{3/2}(x) \quad \text{and} \quad j_2(x) = \sqrt{\frac{\pi}{2x}} J_{5/2}(x).$$

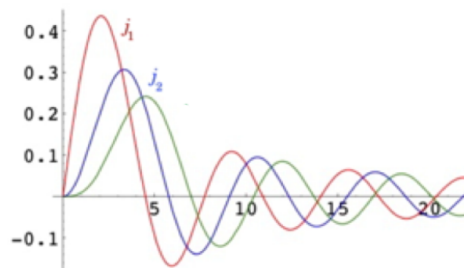
Then from Problem 32 we have

$$J_{3/2}(x) = \sqrt{2\pi x} \left( \frac{\sin x}{x} - \cos x \right) \quad \text{so} \quad j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}$$

and

$$J_{5/2}(x) = \sqrt{2\pi x} \left( \frac{3\sin x}{x^2} - \frac{3\cos x}{x} - \sin x \right) \quad \text{so} \quad j_2(x) = \left( \frac{3}{x^3} - \frac{1}{x} \right) \sin x - \frac{3\cos x}{x^2}$$

**(b)** Using a graphing utility to plot the graphs of  $j_1(x)$  and  $j_2(x)$ , we get the red and blue graphs in the figure to the right.



$$31. f(t) = 2 - 2\mathcal{U}(t-2) + [(t-2) + 2]\mathcal{U}(t-2) = 2 + (t-2)\mathcal{U}(t-2)$$

$$\mathcal{L}\{f(t)\} = \frac{2}{s} + \frac{1}{s^2}e^{-2s}$$

$$\mathcal{L}\{e^t f(t)\} = \frac{2}{s-1} + \frac{1}{(s-1)^2}e^{-2(s-1)}$$

$$32. f(t) = t - t\mathcal{U}(t-1) + (2-t)\mathcal{U}(t-1) - (2-t)\mathcal{U}(t-2) = t - 2(t-1)\mathcal{U}(t-1) + (t-2)\mathcal{U}(t-2)$$

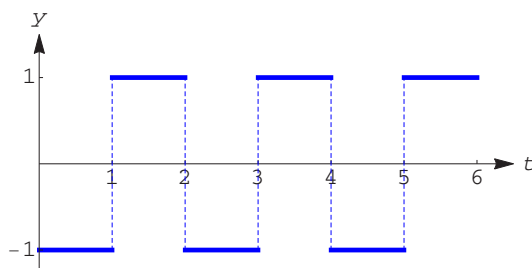
$$\mathcal{L}\{f(t)\} = \frac{1}{s^2} - \frac{2}{s^2}e^{-s} + \frac{1}{s^2}e^{-2s}$$

$$\mathcal{L}\{e^t f(t)\} = \frac{1}{(s-1)^2} - \frac{2}{(s-1)^2}e^{-(s-1)} + \frac{1}{(s-1)^2}e^{-2(s-1)}$$

33. The graph of

$$f(t) = -1 + 2 \sum_{k=1}^{\infty} (-1)^{k+1} \mathcal{U}(t-k) = -1 + 2\mathcal{U}(t-1) - 2\mathcal{U}(t-2) + 2\mathcal{U}(t-3) - \dots$$

is



One way of proceeding to find the Laplace transform is to take the transform term-by-term of the series:

$$\mathcal{L}\{f(t)\} = -\frac{1}{s} + \frac{2}{s}e^{-s} - \frac{2}{s}e^{-2s} + \frac{2}{s}e^{-3s} - \dots \quad \leftarrow \text{geometric series}$$

For  $s > 0$ ,

$$\begin{aligned} \mathcal{L}\{f(t)\} &= -\frac{1}{s} + \frac{2}{s} [e^{-s} - e^{-2s} + e^{-3s} - \dots] = -\frac{1}{s} + \frac{2}{s} \cdot \frac{e^{-s}}{1 + e^{-s}} \\ &= \frac{e^{-s} - 1}{s(1 + e^{-s})} \end{aligned}$$

Alternatively, since  $f$  is a periodic functions it can also be defined by

$$f(t) = \begin{cases} -1, & 0 \leq t < 1 \\ 1, & 1 \leq t < 2, \end{cases} \quad \text{where } f(t+2) = f(t).$$

**I.** If  $\lambda = 0$  then  $X'' = 0$  and  $X(x) = c_1x + c_2$ . Also  $Y'' - Y = 0$  and  $Y(y) = c_3 \cosh y + c_4 \sinh y$  so

$$u = XY = (c_1x + c_2)(c_3 \cosh y + c_4 \sinh y).$$

**II.** If  $\lambda = -\alpha^2 < 0$  then  $X'' - \alpha^2 X = 0$  and  $Y'' + (\alpha^2 - 1)Y = 0$ . The solution of the first differential equation is  $X(x) = c_5 \cosh \alpha x + c_6 \sinh \alpha x$ . The solution of the second differential equation depends on the nature of  $\alpha^2 - 1$ . We consider three cases:

(i) If  $\alpha^2 - 1 = 0$ , or  $\alpha^2 = 1$ , then  $Y(y) = c_7y + c_8$  and

$$u = XY = (c_5 \cosh \alpha x + c_6 \sinh \alpha x)(c_7y + c_8).$$

(ii) If  $\alpha^2 - 1 < 0$ , or  $0 < \alpha^2 < 1$ , then  $Y(y) = c_9 \cosh \sqrt{1 - \alpha^2} y + c_{10} \sinh \sqrt{1 - \alpha^2} y$  and

$$u = XY = (c_5 \cosh \alpha x + c_6 \sinh \alpha x) \left( c_9 \cosh \sqrt{1 - \alpha^2} y + c_{10} \sinh \sqrt{1 - \alpha^2} y \right).$$

(iii) If  $\alpha^2 - 1 > 0$ , or  $\alpha^2 > 1$ , then  $Y(y) = c_{11} \cos \sqrt{\alpha^2 - 1} y + c_{12} \sin \sqrt{\alpha^2 - 1} y$  and

$$u = XY = (c_5 \cosh \alpha x + c_6 \sinh \alpha x) \left( c_{11} \cos \sqrt{\alpha^2 - 1} y + c_{12} \sin \sqrt{\alpha^2 - 1} y \right).$$

**III.** If  $\lambda = \alpha^2 > 0$ , then  $X'' + \alpha^2 X = 0$  and  $X(x) = c_{13} \cos \alpha x + c_{14} \sin \alpha x$ . Also,

$Y'' - (1 + \alpha^2)Y = 0$  and  $Y(y) = c_{15} \cosh \sqrt{1 + \alpha^2} y + c_{16} \sinh \sqrt{1 + \alpha^2} y$  so

$$u = XY = (c_{13} \cos \alpha x + c_{14} \sin \alpha x) \left( c_{15} \cosh \sqrt{1 + \alpha^2} y + c_{16} \sinh \sqrt{1 + \alpha^2} y \right).$$

16. Substituting  $u(x, t) = X(x)T(t)$  into the partial differential equation yields  $a^2 X''T - g = XT''$ , which is not separable.
17. Identifying  $A = B = C = 1$ , we compute  $B^2 - 4AC = -3 < 0$ . The equation is elliptic.
18. Identifying  $A = 3$ ,  $B = 5$ , and  $C = 1$ , we compute  $B^2 - 4AC = 13 > 0$ . The equation is hyperbolic.
19. Identifying  $A = 1$ ,  $B = 6$ , and  $C = 9$ , we compute  $B^2 - 4AC = 0$ . The equation is parabolic.
20. Identifying  $A = 1$ ,  $B = -1$ , and  $C = -3$ , we compute  $B^2 - 4AC = 13 > 0$ . The equation is hyperbolic.
21. Identifying  $A = 1$ ,  $B = -9$ , and  $C = 0$ , we compute  $B^2 - 4AC = 81 > 0$ . The equation is hyperbolic.
22. Identifying  $A = 0$ ,  $B = 1$ , and  $C = 0$ , we compute  $B^2 - 4AC = 1 > 0$ . The equation is hyperbolic.
23. Identifying  $A = 1$ ,  $B = 2$ , and  $C = 1$ , we compute  $B^2 - 4AC = 0$ . The equation is parabolic.



so

$$A_0 = 0, \quad A_1 + B_1 = 0, \quad C_1 + D_1 = 75,$$

and

$$A_n + B_n = 0, \quad C_n + D_n = 0, \quad \text{for } n > 1.$$

When  $r = 2$

$$A_0 + B_0 \ln 2 = \frac{1}{2\pi} \int_0^{2\pi} 60 \cos \theta \, d\theta = 0$$

$$A_n 2^n + B_n 2^{-n} = \frac{1}{\pi} \int_0^{2\pi} 60 \cos \theta \cos n\theta \, d\theta = \begin{cases} 0, & n > 1 \\ 60, & n = 1 \end{cases}$$

$$C_n 2^n + D_n 2^{-n} = \frac{1}{\pi} \int_0^\infty 60 \cos \theta \sin n\theta \, d\theta = 0, \quad n = 1, 2, \dots,$$

so

$$B_0 = 0, \quad 2A_1 + \frac{1}{2}B_1 = 60, \quad 2C_1 + \frac{1}{2}D_1 = 0,$$

and

$$A_n 2^n + B_n 2^{-n} = 0, \quad C_n 2^n + D_n 2^{-n} = 0, \quad \text{for } n > 1.$$

We have  $A_0 = 0$  and  $B_0 = 0$ , and solving the nonhomogeneous systems for  $n = 1$ ,

$$A_1 + B_1 = 0 \qquad C_1 + D_1 = 75$$

$$2A_1 + \frac{1}{2}B_1 = 60 \qquad 2C_1 + \frac{1}{2}D_1 = 0$$

yields  $A_1 = 40$ ,  $B_1 = -40$ ,  $C_1 = -25$ , and  $D_1 = 100$ . Finally, solving the homogeneous systems

$$A_n + B_n = 0 \qquad C_n + D_n = 0$$

$$A_n 2^n + B_n 2^{-n} = 0 \qquad C_n 2^n + D_n 2^{-n} = 0$$

gives  $A_n = B_n = C_n = D_n = 0$  for  $n > 1$ . The solution is then

$$\begin{aligned} u(r, \theta) &= (A_1 r + B_1 r^{-1}) \cos \theta + (C_1 r + D_1 r^{-1}) \sin \theta \\ &= (4 - r - 40r^{-1}) \cos \theta + (-25r + 100r^{-1}) \sin \theta \\ &= 40 \left( r - \frac{1}{r} \right) \cos \theta - 25 \left( r - \frac{4}{r} \right) \sin \theta. \end{aligned}$$

14. We solve

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0, \quad 0 < \theta < \pi, \quad a < r < b,$$

$$u(a, \theta) = \theta(\pi - \theta), \quad u(b, \theta) = 0, \quad 0 < \theta < \pi,$$

$$u(r, 0) = 0, \quad u(r, \pi) = 0, \quad a < r < b.$$

By writing the boundary condition  $x = 0$  as

$$u(0, t) = u_0 - u_0 \mathcal{U}(t - 1)$$

its transform is

$$\begin{aligned} U(0, s) &= \frac{u_0}{s} - \frac{u_0}{s} e^{-s} \\ c_1 &= \frac{u_0}{s} - \frac{u_0}{s} e^{-s} \\ U(x, s) &= u_0 \frac{e^{-\sqrt{s}x}}{s} - u_0 \frac{e^{-\sqrt{s}x}}{s} e^{-s} \\ u(x, t) &= u_0 \mathcal{L}^{-1} \left\{ \frac{e^{-\sqrt{s}x}}{s} \right\} - u_0 \mathcal{L}^{-1} \left\{ \frac{e^{-\sqrt{s}x}}{s} e^{-s} \right\} \end{aligned}$$

by entry 3 of Table 14.1.1 and the inverse form of the second translation theorem that:

$$u(x, t) = u_0 \operatorname{erfc} \left( \frac{x}{2\sqrt{t}} \right) - u_0 \operatorname{erfc} \left( \frac{x}{2\sqrt{t-1}} \right) \mathcal{U}(t - 1)$$

or

$$u(x, t) = \begin{cases} u_0 \operatorname{erfc} \left( \frac{x}{2\sqrt{t}} \right), & 0 < t < 1 \\ u_0 \operatorname{erfc} \left( \frac{x}{2\sqrt{t}} \right) - u_0 \operatorname{erfc} \left( \frac{x}{2\sqrt{t-1}} \right), & t > 1. \end{cases}$$

**18.** The Laplace transform with respect to  $t$  of the partial differential equation gives

$$\frac{d^2 U}{dx^2} - sU = -50 \quad \text{so} \quad U(x, s) = c_1 e^{-\sqrt{s}x} + c_2 e^{\sqrt{s}x} + \frac{50}{s}.$$

The boundary condition

$$\lim_{x \rightarrow \infty} u(x, t) = 50 \quad \text{implies} \quad \lim_{x \rightarrow \infty} U(x, s) = \frac{50}{s}$$

so we take  $c_2 = 0$ . Thus

$$U(x, s) = c_1 e^{-\sqrt{s}x} + \frac{50}{s}.$$

The transform of the boundary condition at  $x = 0$  is

$$U(0, s) = \frac{100}{s} e^{-5s} - \frac{100}{s} e^{-10s}.$$

Since

$$\frac{100}{s} e^{-5s} - \frac{100}{s} e^{-10s} = c_1 + \frac{50}{s}$$

we have

$$c_1 = -\frac{50}{s} + \frac{100}{s} e^{-5s} - \frac{100}{s} e^{-10s}$$